

Physics 452: Homework #20

Due Thursday, Apr. 2, 5:00PM, 2009

1. If $\bar{A}(\bar{r}, t) = \hat{z}A_0 \sin(kx - \omega t)$ and $\phi(\bar{r}) = \frac{e}{4\pi\epsilon_0 r}$, compute \bar{E} and \bar{B} .

2. By installing $\bar{A}(\bar{r}, t) = A_x(\bar{r}, t)\hat{x} + A_y(\bar{r}, t)\hat{y} + A_z(\bar{r}, t)\hat{z}$ show explicitly that

$$\bar{F}_{Lorentz} = q \left[-\nabla\phi - \frac{\partial\bar{A}}{\partial t} + \bar{v} \times (\nabla \times \bar{A}) \right] \text{ and}$$

$$\bar{F}_{Lorentz} = \nabla \left(-mc^2 \sqrt{1 - v^2/c^2} - q\phi + q\bar{v} \cdot \bar{A} \right) - \frac{d}{dt} \nabla_{\bar{v}} (-q\phi + q\bar{v} \cdot \bar{A}) \text{ are equivalent.}$$

HINT: Straight away, you can write $\nabla \left(\sqrt{1 - v^2/c^2} \right) = 0$, $\nabla_{\bar{v}}(\phi) = 0$, and $\nabla_{\bar{v}}(\bar{v} \cdot \bar{A}) = \bar{A}$.

Note also that $\frac{\partial\bar{A}}{\partial t} \neq \frac{d\bar{A}}{dt}$.

3. Nonrelativistically, we have $\bar{F} = \frac{d\bar{p}}{dt} = \frac{dm\bar{v}}{dt} = \frac{d}{dt} \nabla_{\bar{v}} \left(\frac{1}{2} m v^2 \right)$. In this case, determine the

Lagrangian, canonical momentum, and Hamiltonian for a charged particle experiencing electromagnetic potentials ϕ and \bar{A} . Show that quantization leads to the following form

of the Schrödinger equation: $i\hbar \frac{\partial\Psi}{\partial t} = \frac{1}{2m} (-i\hbar\nabla - q\bar{A})^2 \Psi + q\phi\Psi$.

4. For the fields in problem 1, justify our use of

$$\Rightarrow i\hbar \frac{\partial\Psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \Psi + q \frac{e}{4\pi\epsilon_0 r} \Psi + q\omega A_0 z \cos(\omega t) \Psi \text{ in chapter 9 of the textbook, starting from}$$

$$i\hbar \frac{\partial\Psi}{\partial t} = \frac{1}{2m} (-i\hbar\nabla - q\bar{A})^2 \Psi + q\phi\Psi.$$

HINT: Note: $(-i\hbar\nabla - q\bar{A})^2 \Psi = (-\hbar^2\nabla^2 + i\hbar q\nabla \cdot \bar{A} + i\hbar q\bar{A} \cdot \nabla + q^2 A^2) \Psi$ where

$\nabla \cdot \bar{A}\Psi = (\nabla \cdot \bar{A})\Psi + \bar{A} \cdot \nabla\Psi$. Also, $\nabla \cdot \bar{A} = 0$ for our field. Next let $\Psi(\bar{r}, t) \equiv \tilde{\Psi}(\bar{r}, t) e^{-i\frac{q}{\hbar}\bar{A} \cdot \bar{r}}$, and obtain an equation for $\tilde{\Psi}$. You may also make the dipole approximation $kx \ll 1$.